# On stability in the Borg-Hochstadt theorem for periodic Jacobi matrices 

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To V. A. Marchenko on occasion of his $95^{\text {th }}$ anniversary


#### Abstract

A result of Borg-Hochstadt in the theory of periodic Jacobi matrices states that such a matrix has constant diagonals as long as all gaps in its spectrum are closed (have zero length). We suggest a quantitative version of this result by proving two-sided bounds between oscillations of the matrix entries along the diagonals and the length of the maximal gap in the spectrum.


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## 1. Introduction

In his recent paper [9] V. A. Marchenko reverts to the classical subject of periodic Jacobi matrices. He gives an intrinsic description of polynomials with all their $\pm 2$-points real and so obtains parametrization of the Hill discriminants of such matrices. His argument is straightforward and makes no appeal to conformal mappings.

The Hill discriminants play a crucial role in the problem related to periodic Jacobi matrices we address in this note. Recall that Jacobi matrices are two-sided, infinite, three-diagonal matrices of the form

$$
J=\left[\begin{array}{ccccccc}
\ddots & \ddots & \ddots & & & &  \tag{1.1}\\
& a_{-1} & b_{0} & a_{0} & & & \\
& & a_{0} & b_{1} & a_{1} & & \\
& & & a_{1} & b_{2} & a_{2} & \\
& & & & \ddots & \ddots & \ddots
\end{array}\right], \quad b_{j} \in \mathbb{R}, a_{j}>0
$$

$J$ is said to be periodic of period $p \in \mathbb{N}:=\{1,2, \ldots\}$, if

$$
a_{j+p}=a_{j}, \quad b_{j+p}=b_{j}, \quad j \in \mathbb{Z}:=\{0, \pm 1, \pm 2, \ldots\}
$$

A p-periodic Jacobi matrix $J$ (1.1) generates in an obvious way a bounded, self-adjoint, linear operator $J$ on the Hilbert space $\ell^{2}(\mathbb{Z})$. Its spectrum $\sigma(J)$ is known to have a banded structure, i.e., it is composed of $p$ spectral bands (closed intervals)

$$
\begin{equation*}
\sigma(J)=\bigcup_{j=0}^{p-1}\left[\mu_{j}^{+}, \mu_{j+1}^{-}\right], \quad \mu_{0}^{+}<\mu_{1}^{-} \leq \mu_{1}^{+}<\cdots<\mu_{p-1}^{-} \leq \mu_{p-1}^{+}<\mu_{p}^{-} \tag{1.2}
\end{equation*}
$$

some of them can merge. A convex hall of the spectrum (the least interval that contains the whole spectrum) is $L=\left[\mu_{0}^{+}, \mu_{p}^{-}\right]$.

The bands are interspersed with (interior) gaps

$$
\begin{equation*}
\gamma_{j}:=\left(\mu_{j}^{-}, \mu_{j}^{+}\right), \quad j=1,2, \ldots, p-1, \quad \mu_{j}^{-} \leq \mu_{j}^{+} \tag{1.3}
\end{equation*}
$$

of the length $\left|\gamma_{j}\right|=\mu_{j}^{+}-\mu_{j}^{-}$, and $\mu_{j}^{-}=\mu_{j}^{+}$means that the gap is closed (the adjacent bands merge). Yet it seems advisable viewing a closed gap as an actual gap (of zero length) rather than dealing with two merging bands as a single one. We stick to this viewpoint consistently throughout the paper. We observe the situation, when the closed gaps arise, in the simplest example of constant Jacobi matrices $J_{0}$ with $a_{j}=a_{0}, b_{j}=b_{0}, j \in \mathbb{Z}$. Now all gaps are closed, and the spectrum $\sigma\left(J_{0}\right)=\left[b_{0}-2 a_{0}, b_{0}+2 a_{0}\right]$ is a single interval. Following the above convention, we can (and will) view such matrices as periodic of period $p \in \mathbb{N}$.

The well-known result of Borg-Hochstadt [3, 4], (see [11, Theorem 5.4.21]), states, that the converse is also true. Precisely, a periodic Jacobi matrix $J$ with all gaps closed is constant, $J=J_{0}$. For a simple linear algebraic approach to this result see [2]. The problem we address here is stability (or a quantitative
version) of this result. Specifically, we show that for periodic Jacobi matrices with a "small" variation of parameters $a$ 's and $b$ 's, the gaps in their spectra are "small", and vice versa. The quantitative stability in the original Borg's theorem for the Hill operators on the real line was studied in [5].

To be more formal, given a bounded sequence $c=\left\{c_{j}\right\}_{j \in \mathbb{Z}}$ of real numbers, its variation is defined by

$$
\begin{equation*}
\omega_{c}:=\sup _{i, j \in \mathbb{Z}}\left(c_{i}-c_{j}\right) \geq 0 \tag{1.4}
\end{equation*}
$$

Let $\gamma$ be a maximal gap in the spectrum of $J$, so

$$
\begin{equation*}
|\gamma|=\max _{1 \leq j \leq p-1}\left|\gamma_{j}\right| \tag{1.5}
\end{equation*}
$$

Here is the main result of the paper.
Theorem 1.1. Let $J$ be a periodic Jacobi matrix (1.1) of period $p$. Then

$$
\begin{equation*}
\omega_{b} \leq(p-1)|\gamma|, \quad \omega_{a} \leq p^{2} \sqrt{p}|\gamma| \tag{1.6}
\end{equation*}
$$

Furthermore,

$$
\begin{equation*}
\omega_{b}+\omega_{a} \geq \frac{|\gamma|}{4} \tag{1.7}
\end{equation*}
$$

In particular, $|\gamma|=0$ (all gaps are closed) if and only if $\omega_{a}=\omega_{b}=0(J$ is a constant Jacobi matrix), so the result of Borg-Hochstadt follows. In the case of periodic Jacobi matrices the supremum in the definition (1.4) is obviously attained, for the matrix entries take only a finite number of values.

A periodic Jacobi matrix is said to be normalized if

$$
\begin{equation*}
\sum_{j=1}^{p} b_{j}=0 \tag{1.8}
\end{equation*}
$$

Clearly, each Jacobi matrix can be normalized by adding an appropriate constant to all b's. Under such transformation (shift) neither the variation of the entries along diagonals, nor the length of the maximal gap alters, so throughout the paper we assume (1.8) to hold.

As in [9], the argument is by and large elementary, and relies upon some basic properties of the Hill discriminant with regard to the spectrum $\sigma(J)$, see [10, Chapter 7], [11, Chapter 5].

A key ingredient of the proof is the extremal problem "à la" the Chebyshev Alternance Theorem, suggested by Korotyaev and Kutsenko [7, Lemma 2.2].

Let $x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}, n \geq 2$, and put

$$
X(\lambda)=X(\lambda, x):=\prod_{j=1}^{n}\left(\lambda-x_{j}\right)
$$

A set $\mathcal{P}_{n}(c) \subset \mathbb{R}^{n}$ is defined by imposing the following conditions: $x=\left(x_{1}, \ldots, x_{n}\right) \in \mathcal{P}_{n}(c)$ if and only if
(1) $x_{1}+\cdots+x_{n}=0$;
(2) $|X(\lambda)| \leq c, \quad \lambda \in\left[\min _{j} x_{j}, \max _{j} x_{j}\right]$.

The result of Korotyaev-Kutsenko states that

$$
\begin{equation*}
\sup \left\{\|x\|^{2}=x_{1}^{2}+\cdots+x_{n}^{2}: x \in \mathcal{P}_{n}(c)\right\}=2 n\left(\frac{c}{2}\right)^{2 / n} \tag{1.9}
\end{equation*}
$$

It is instructive to paraphrase this result in a quantitative form. An algebraic polynomial $P$ with the roots $\left\{x_{j}\right\}$ is balanced, if $x_{1}+\cdots+x_{n}=0$.

Theorem KK. Let

$$
P(x)=\prod_{j=1}^{n}\left(x-x_{j}\right)=x^{n}+\sigma_{2} x^{n-2}+\cdots, \quad x_{1} \leq \cdots \leq x_{n}
$$

be a balanced polynomial with the real roots $\left\{x_{j}\right\}$. Then

$$
\begin{equation*}
\sum_{j=1}^{n} x_{j}^{2} \leq 2 n\left(\frac{\|P\|}{2}\right)^{2 / n}, \quad\|P\|:=\max _{x \in\left[x_{1}, x_{n}\right]}|P(x)| \tag{1.10}
\end{equation*}
$$

Clearly, the equality prevails in (1.10) for the monic Chebyshev polynomials of the first kind.

The work on this note was inspired to a large extent by a recent paper [8]. Up to distinction in terminology and notation, [8] contains the upper bound for $\omega_{b}$ and the lower bound (1.7). The main contribution of this note is a new upper bound for $\omega_{a}$ in (1.6), providing thereby the quantitative stability in the Borg-Hochstadt theorem in its full generality.

We illustrate our results on the example of 4-periodic Jacobi matrices.

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## 2. Proof of the main result: upper bounds

We recall some rudiments of the theory of Jacobi matrices [10, 11].
The basic recurrence relation for Jacobi matrices (1.1) is of the form

$$
\begin{equation*}
J y=\lambda y \sim a_{n-1} y_{n-1}+b_{n} y_{n}+a_{n} y_{n+1}=\lambda y_{n}, \quad \lambda \in \mathbb{C}, n \in \mathbb{Z}, \tag{2.1}
\end{equation*}
$$

and a pair of special solutions of (2.1), $\left\{c_{n}(\cdot, m)\right\}_{n \in \mathbb{Z}}$ and $\left\{s_{n}(\cdot, m)\right\}_{n \in \mathbb{Z}}, m \in \mathbb{Z}$, with the initial data

$$
\begin{array}{ll}
s_{m-1}(\lambda, m)=0, & s_{m}(\lambda, m)=1, \\
c_{m-1}(\lambda, m)=-1, & c_{m}(\lambda, m)=0, \tag{2.2}
\end{array}
$$

is of particular concern. It is clear that $s_{m+j}(\cdot, m)\left(c_{m+j}(\cdot, m)\right)$ is the algebraic polynomial of degree $j(j-1)$, respectively.

The polynomial $s_{m+p-1}$ is crucial for the theory of periodic Jacobi matrices. Its zeros $\left\{\xi_{j}^{(m)}\right\}_{j=1}^{p-1}$ are known to be simple and real, and

$$
\begin{align*}
a^{p} s_{m+p-1}(\lambda, m) & =\prod_{j=1}^{p-1}\left(\lambda-\xi_{j}^{(m)}\right)=\lambda^{p-1}-\lambda^{p-2} \sum_{j=0}^{p-2} b_{m+j}+\cdots,  \tag{2.3}\\
a & :=\left(a_{1} \ldots a_{p}\right)^{1 / p} .
\end{align*}
$$

In view of normalization (1.8) and periodicity,

$$
\sum_{j=1}^{p-1} \xi_{j}^{(m)}=\sum_{j=0}^{p-2} b_{m+j}=-b_{m+p-1}=-b_{m-1}
$$

so for each $k, l \in \mathbb{Z}$ the following trace formula is valid

$$
\begin{equation*}
b_{k-1}-b_{l-1}=\sum_{j=1}^{p-1}\left(\xi_{j}^{(l)}-\xi_{j}^{(k)}\right) . \tag{2.4}
\end{equation*}
$$

On the other hand, for each $m \in \mathbb{Z}$ the zero $\xi_{j}^{(m)} \in \bar{\gamma}_{j}=\left[\mu_{j}^{-}, \mu_{j}^{+}\right]$, see, e.g., [10, formula (7.50)] and [11, Theorem 5.4.16], so

$$
\left|b_{k-1}-b_{l-1}\right|=\sum_{j=1}^{p-1}\left(\xi_{j}^{(l)}-\xi_{j}^{(k)}\right) \leq(p-1)|\gamma|,
$$

and the upper bound for $\omega_{b}$ in (1.6) follows.

The upper bound for $\omega_{a}$ is much more intricate. Now the Hill discriminant $\mathcal{D}$ appears on the central stage (see [11, Section 5.4] for a detailed account of discriminants and their properties). This object arises in two guises: as a difference of the special solutions of (2.1) (or, equivalently, the trace of the monodromy matrix)

$$
\mathcal{D}(\lambda)=s_{m+p}(\lambda, m)-c_{m+p-1}(\lambda, m),
$$

and as a characteristic polynomial of a certain Hermitian matrix

$$
a^{p} \mathcal{D}(\lambda)=\operatorname{det}\left(\lambda-\Phi_{m}\right)=\prod_{j=1}^{p}\left(\lambda-d_{j}\right)
$$

where $a$ is defined in (2.3),

$$
\Phi_{m}=\left[\begin{array}{cccccc}
b_{m} & a_{m} & & & & a_{m+p-1} i  \tag{2.5}\\
a_{m} & b_{m+1} & a_{m+1} & & & \\
& a_{m+1} & b_{m+2} & a_{m+2} & & \\
& & \ddots & \ddots & \ddots & \\
& & & a_{m+p-2} & b_{m+p-2} & a_{m+p-2} \\
-a_{m+p-1} i & & & & a_{m+p-1} & b_{m+p-1}
\end{array}\right]
$$

As a matter of fact, $\mathcal{D}$ does not depend on $m$, and neither do its zeros $\left\{d_{j}\right\}$. By normalization (1.8) and (2.5), $\mathcal{D}$ is the balanced polynomial of degree $p$.

The zeros $\left\{d_{j}\right\}_{j=1}^{p}$ of $\mathcal{D}$ (or, equivalently, the eigenvalues of $\Phi_{m}$ ) are known to be real and simple. Evaluating of the trace of $\Phi_{m}^{2}$ in two ways provides, in view of periodicity, another trace formula

$$
\begin{equation*}
\operatorname{tr} \Phi_{m}^{2}=\sum_{j=1}^{p} d_{j}^{2}=\sum_{j=1}^{p}\left(b_{j}^{2}+2 a_{j}^{2}\right) \tag{2.6}
\end{equation*}
$$

There is a tight relation between the Hill discriminant and the spectrum of the underlying Jacobi matrix $J$. Precisely,

$$
\sigma(J)=\mathcal{D}^{(-1)}[-2,2], \quad\left|\mathcal{D}\left(\mu_{j}^{ \pm}\right)\right|=2
$$

so $\sigma(J)$ is the inverse image of the interval $[-2,2]$, and $|\mathcal{D}|=2$ at all endpoints of the spectrum. In other words, the set $\left\{\mu_{j}^{ \pm}\right\}_{j}$ constitutes the full collection of $\pm 2$-points of the Hill discriminant.

Denote by $2 M$ the sup-norm of the discriminant over the convex hall of the spectrum

$$
2 M=\|\mathcal{D}\|=\|\mathcal{D}\|_{C(L)}, \quad L=\left[\mu_{0}^{+}, \mu_{p}^{-}\right]
$$

Then, $M \geq 1$, and by the Borg-Hochstadt theorem, $M=1$ if and only if $|\gamma|=0$, so $J$ is the constant matrix. We exclude this case in what follows, so let $M>1$. The discriminant attains this extremum at some critical point $v_{k} \in \gamma_{k}=\left(\mu_{k}^{-}, \mu_{k}^{+}\right)$. By the Taylor formula, there is $\xi \in\left[\mu_{k}^{-}, \mu_{k}^{+}\right]$so that for $\mu_{k}^{-} \leq t \leq \mu_{k}^{+}$the equality

$$
\begin{aligned}
\mathcal{D}(t) & =\mathcal{D}\left(v_{k}\right)+\mathcal{D}^{\prime}\left(v_{k}\right)\left(t-v_{k}\right)+\frac{\mathcal{D}^{\prime \prime}(\xi)}{2}\left(t-v_{k}\right)^{2} \\
& =\mathcal{D}\left(v_{k}\right)+\frac{\mathcal{D}^{\prime \prime}(\xi)}{2}\left(t-v_{k}\right)^{2}
\end{aligned}
$$

holds. Putting $t=\mu_{k}^{+}$provides

$$
2(M-1)=\frac{\left|\mathcal{D}^{\prime \prime}(\xi)\right|}{2}\left(\mu_{k}^{+}-v_{k}\right)^{2}
$$

or

$$
\left(\mu_{k}^{+}-v_{k}\right)^{2}=\frac{4(M-1)}{\left|\mathcal{D}^{\prime \prime}(\xi)\right|} \geq \frac{4(M-1)}{\left\|\mathcal{D}^{\prime \prime}\right\|}
$$

By A. Markov's inequality for the interval $L$ [1, Theorem 5.1.8],

$$
\left\|\mathcal{D}^{\prime \prime}\right\| \leq \frac{4 p^{4} M}{|L|^{2}}
$$

and we end up with the lower bound for the length of the maximal gap

$$
\begin{equation*}
|\gamma|^{2} \geq\left(\mu_{k}^{+}-v_{k}\right)^{2} \geq \frac{|L|^{2}}{p^{4}} \cdot \frac{M-1}{M} \Longrightarrow|\gamma| \geq \frac{|L|}{p^{2}} \cdot\left(\frac{M-1}{M}\right)^{1 / 2} \tag{2.7}
\end{equation*}
$$

We distinguish two cases.

1. The case when $M \geq 2$ is simple. We show that not only the variation of $a_{j}$ 's, but the numbers $a_{j}$ 's themselves are "small." Indeed, it follows now from (2.7) that

$$
\begin{equation*}
|\gamma| \geq \frac{|L|}{\sqrt{2} p^{2}} \tag{2.8}
\end{equation*}
$$

Note that for an algebraic polynomial of degree $p \geq 2$ the sum of its $\alpha$-points does not depend on $\alpha$. In particular, for a balanced polynomial the sum of its $\alpha$-points is zero. Hence, by (1.8),

$$
\mu_{0}^{+}+\mu_{p}^{-}+\sum_{j=1}^{p-1}\left(\mu_{j}^{-}+\mu_{j}^{+}\right)=2 \sum_{j=1}^{p} b_{j}=0
$$

(the left-hand side being the sum of all $\pm 2$-points of $\mathcal{D}$ ), so $\mu_{0}^{+}<0<\mu_{p}^{-}$. Since $d_{j} \in \sigma(J)$, we have

$$
\left|d_{j}\right| \leq|L|, \quad j=1,2, \ldots, p
$$

and the trace formula (2.6) implies

$$
\sum_{j=1}^{p}\left(b_{j}^{2}+2 a_{j}^{2}\right)=\sum_{j=1}^{p} d_{j}^{2}<p|L|^{2}
$$

In view of (2.8), we come to the bound

$$
\begin{equation*}
\max _{j} a_{j}<\sqrt{\frac{p}{2}}|L| \leq p^{2} \sqrt{p}|\gamma| \tag{2.9}
\end{equation*}
$$

as claimed.
2. Let us turn to the harder part $1<M<2$. The extremal problem of Korotyaev-Kutsenko comes into play here. We apply Theorem KK to the balanced polynomial $P=a^{p} \mathcal{D}, a$ is defined in (2.3), to obtain

$$
\sum_{j=1}^{p} d_{j}^{2} \leq 2 p\left(\frac{1}{2}\|P\|_{C\left[d_{1}, d_{p}\right]}\right)^{2 / p} \leq 2 p a^{2} M^{2 / p}
$$

By the trace formula (2.6),

$$
\frac{1}{p} \sum_{j=1}^{p} a_{j}^{2} \leq a^{2} M^{2 / p}
$$

and the arithmetic-geometric means (AGM for short) inequality yields

$$
\begin{equation*}
\frac{1}{p} \sum_{j=1}^{p} a_{j}^{2}-\left(\frac{1}{p} \sum_{j=1}^{p} a_{j}\right)^{2} \leq \frac{1}{p} \sum_{j=1}^{p} a_{j}^{2}-a^{2} \leq a^{2}\left(M^{2 / p}-1\right) \tag{2.10}
\end{equation*}
$$

But the left-hand side of (2.10) is

$$
\frac{1}{p} \sum_{j=1}^{p} a_{j}^{2}-\left(\frac{1}{p} \sum_{j=1}^{p} a_{j}\right)^{2}=\frac{1}{p} \sum_{j=1}^{p}\left(a_{j}-s_{a}\right)^{2}, \quad s_{a}:=\frac{1}{p} \sum_{j=1}^{p} a_{j}
$$

To estimate the right-hand side of (2.10) from above, we apply the well-known inequality

$$
u^{\tau}-1 \leq \tau(u-1), \quad u>0,0<\tau \leq 1
$$

Thus, we come to the upper bound

$$
\frac{1}{p} \sum_{j=1}^{p}\left(a_{j}-s_{a}\right)^{2} \leq \frac{2 a^{2}}{p}(M-1)
$$

To proceed further, recall the assumption $1<M<2$. From (2.7) we derive

$$
|\gamma| \geq \frac{|L|}{\sqrt{2} p^{2}} \sqrt{M-1}, \quad M-1 \leq \frac{2 p^{4}}{|L|^{2}}|\gamma|^{2},
$$

and so

$$
\frac{1}{p} \sum_{j=1}^{p}\left(a_{j}-s_{a}\right)^{2} \leq \frac{4 a^{2} p^{3}}{|L|^{2}}|\gamma|^{2} .
$$

By [6, Lemma 2.2], $|L|=\mu_{p}^{-}-\mu_{0}^{+} \geq 4 a$, so, finally,

$$
\begin{equation*}
\frac{1}{p} \sum_{j=1}^{p}\left(a_{j}-s_{a}\right)^{2} \leq \frac{p^{3}}{4}|\gamma|^{2}, \quad\left|a_{j}-s_{a}\right| \leq \frac{p^{2}}{2}|\gamma| \tag{2.11}
\end{equation*}
$$

for each $j=1,2, \ldots, p$. The latter inequality obviously implies (1.6) for $\omega_{a}$. The proof of the upper bounds in Theorem 1.1 is complete.

## 3. 4-periodic Jacobi matrices

Consider an instructive example of 4-periodic Jacobi matrices $J$ with $b_{j} \equiv 0$, $j \in \mathbb{Z}$. The Hill discriminant is now

$$
\begin{gathered}
\mathcal{D}(\lambda)=s_{m+4}(\lambda, m)-c_{m+3}(\lambda, m)=a^{-4}\left(\lambda^{4}-\alpha \lambda^{2}+\beta\right), \\
a^{4}=a_{1} a_{2} a_{3} a_{4}, \quad \alpha:=a_{1}^{2}+a_{2}^{2}+a_{3}^{2}+a_{4}^{2}, \quad \beta:=\left(a_{1} a_{3}\right)^{2}+\left(a_{2} a_{4}\right)^{2} .
\end{gathered}
$$

The equations $\mathcal{D}(\lambda) \pm 2=0$ are biquadratic, and their roots are

$$
\begin{aligned}
& \mathcal{D}(\lambda)-2=0 \Longrightarrow \lambda=\left\{ \pm \lambda_{1}^{-}, \pm \lambda_{1}^{+}\right\}, \\
& \mathcal{D}(\lambda)+2=0 \Longrightarrow \lambda=\left\{ \pm \lambda_{2}^{-}, \pm \lambda_{2}^{+}\right\} .
\end{aligned}
$$

The above biquadratic equations can be solved explicitly

$$
\lambda_{1}^{ \pm}=\sqrt{\frac{\alpha \pm \sqrt{D^{+}}}{2}}, \quad \lambda_{2}^{ \pm}=\sqrt{\frac{\alpha \pm \sqrt{D^{-}}}{2}},
$$

where

$$
\begin{aligned}
D^{+} & :=\alpha^{2}-4\left(\beta-2 a^{4}\right)=\alpha^{2}-4\left(a_{1} a_{3}-a_{2} a_{4}\right)^{2} \\
& =\left\{\left(a_{1}-a_{3}\right)^{2}+\left(a_{2}+a_{4}\right)^{2}\right\}\left\{\left(a_{1}+a_{3}\right)^{2}+\left(a_{2}-a_{4}\right)^{2}\right\}>0, \\
D^{-} & :=\alpha^{2}-4\left(\beta+2 a^{4}\right)=\alpha^{2}-4\left(a_{1} a_{3}-a_{2} a_{4}\right)^{2} \\
& =\left\{\left(a_{1}-a_{3}\right)^{2}+\left(a_{2}-a_{4}\right)^{2}\right\}\left\{\left(a_{1}+a_{3}\right)^{2}+\left(a_{2}+a_{4}\right)^{2}\right\}>0 .
\end{aligned}
$$

By the AGM inequality,

$$
\begin{equation*}
\sqrt{D^{ \pm}} \leq \alpha \tag{3.1}
\end{equation*}
$$

so all roots are real.
The spectrum is composed of four spectral bands symmetric with respect to the origin

$$
\operatorname{Sp}(J)=\left[\lambda_{1}^{-}, \lambda_{2}^{-}\right] \cup\left[\lambda_{2}^{+}, \lambda_{1}^{+}\right] \cup\left[-\lambda_{2}^{-},-\lambda_{1}^{-}\right] \cup\left[-\lambda_{1}^{+},-\lambda_{2}^{+}\right],
$$

with $0 \leq \lambda_{1}^{-}<\lambda_{2}^{-} \leq \lambda_{2}^{+}<\lambda_{1}^{+}$. The bands are separated with three gaps, the interior one $\gamma_{\mathrm{int}}=\left(-\lambda_{1}^{-}, \lambda_{1}^{-}\right)$, and two symmetric exterior ones, $\pm \gamma_{\mathrm{ext}}$, $\gamma_{\mathrm{ext}}=\left(\lambda_{2}^{-}, \lambda_{2}^{+}\right)$. In view of (3.1), the length of $\gamma_{\mathrm{ext}}$ is bounded from below by

$$
\begin{aligned}
\sqrt{2}\left|\gamma_{\mathrm{ext}}\right| & =\sqrt{2}\left(\lambda_{2}^{+}-\lambda_{2}^{-}\right) \\
& =\sqrt{\alpha+\sqrt{D^{-}}}-\sqrt{\alpha-\sqrt{D^{-}}} \\
& =\frac{2 \sqrt{D^{-}}}{\sqrt{\alpha+\sqrt{D^{-}}}+\sqrt{\alpha-\sqrt{D^{-}}}} \\
& \geq \sqrt{\frac{D^{-}}{2 \alpha}}
\end{aligned}
$$

Therefore,

$$
\begin{aligned}
2\left|\gamma_{\mathrm{ext}}\right| & \geq \sqrt{\frac{D^{-}}{\alpha}} \\
& =\sqrt{\left(a_{1}-a_{3}\right)^{2}+\left(a_{2}-a_{4}\right)^{2}} \sqrt{\frac{\left(a_{1}+a_{3}\right)^{2}+\left(a_{2}+a_{4}\right)^{2}}{\alpha}} \\
& \geq \sqrt{\left(a_{1}-a_{3}\right)^{2}+\left(a_{2}-a_{4}\right)^{2}}
\end{aligned}
$$

(see the definition of $\alpha$ ), and so

$$
\begin{equation*}
\max \left(\left|a_{1}-a_{3}\right|,\left|a_{2}-a_{4}\right|\right) \leq 2\left|\gamma_{\mathrm{ext}}\right| . \tag{3.2}
\end{equation*}
$$

In particular, if the exterior gaps are closed, then $a_{1}=a_{3}, a_{2}=a_{4}$, and the actual period is 2 .

As far as the interior gap goes, we see that

$$
\frac{1}{2}\left|\gamma_{\mathrm{int}}\right|=\lambda_{1}^{-}=\sqrt{\frac{\alpha-\sqrt{D^{+}}}{2}}=\sqrt{2} \frac{\left|a_{1} a_{3}-a_{2} a_{4}\right|}{\sqrt{\alpha+\sqrt{D^{+}}}} \geq \frac{\left|a_{1} a_{3}-a_{2} a_{4}\right|}{\sqrt{\alpha}}
$$

Note that if the interior gap is closed, then $a_{1} a_{3}=a_{2} a_{4}$ (the period may still be 4 ).

Assume with no loss of generality, that $a_{1}+a_{2} \geq a_{3}+a_{4}$. Since

$$
1 \leq \frac{a_{1}+a_{2}+a_{3}+a_{4}}{\sqrt{\alpha}} \leq 2
$$

we come to the bounds

$$
\begin{equation*}
\frac{1}{2} \leq \frac{a_{1}+a_{2}+a_{3}+a_{4}}{2 \sqrt{\alpha}} \leq \frac{a_{1}+a_{2}}{\sqrt{\alpha}} \leq \frac{a_{1}+a_{2}+a_{3}+a_{4}}{\sqrt{\alpha}} \leq 2 \tag{3.3}
\end{equation*}
$$

By (3.3),

$$
\frac{\left|a_{1}-a_{2}\right|}{2} \leq \frac{a_{1}+a_{2}}{\sqrt{\alpha}}\left|a_{1}-a_{2}\right|=\frac{\left|a_{1}^{2}-a_{2}^{2}\right|}{\sqrt{\alpha}}
$$

On the other hand, we write

$$
a_{1}^{2}-a_{2}^{2}=\left(a_{1} a_{3}-a_{2} a_{4}\right)+a_{1}\left(a_{1}-a_{3}\right)+a_{2}\left(a_{4}-a_{2}\right)
$$

to obtain, in view of (3.2), (3.3),

$$
\begin{aligned}
\frac{\left|a_{1}^{2}-a_{2}^{2}\right|}{\sqrt{\alpha}} & \leq \frac{\left|a_{1} a_{3}-a_{2} a_{4}\right|}{\sqrt{\alpha}}+\frac{a_{1}}{\sqrt{\alpha}}\left|a_{1}-a_{3}\right|+\frac{a_{2}}{\sqrt{\alpha}}\left|a_{2}-a_{4}\right| \\
& \leq \frac{1}{2}\left|\gamma_{\mathrm{int}}\right|+2\left|\gamma_{\mathrm{ext}}\right| \frac{a_{1}+a_{2}}{\sqrt{\alpha}} \\
& \leq \frac{1}{2}\left|\gamma_{\mathrm{int}}\right|+4\left|\gamma_{\mathrm{ext}}\right|
\end{aligned}
$$

Hence,

$$
\left|a_{1}-a_{2}\right| \leq\left|\gamma_{\mathrm{int}}\right|+8\left|\gamma_{\mathrm{ext}}\right| \leq 9|\gamma|, \quad|\gamma|:=\max \left(\left|\gamma_{\mathrm{int}}\right|,\left|\gamma_{\mathrm{ext}}\right|\right)
$$

is the length of the maximal gap in the spectrum.
Next, by (3.2),

$$
\left|a_{3}-a_{4}\right| \leq\left|a_{3}-a_{1}\right|+\left|a_{1}-a_{2}\right|+\left|a_{2}-a_{4}\right| \leq\left|\gamma_{\mathrm{int}}\right|+12\left|\gamma_{\mathrm{ext}}\right|,
$$

and we end up with the following bound for the variation of the a's diagonal

$$
\begin{equation*}
\omega_{a} \leq\left|\gamma_{\mathrm{int}}\right|+12\left|\gamma_{\mathrm{ext}}\right| \leq 13|\gamma| \tag{3.4}
\end{equation*}
$$

The opposite inequalities can be proved along the same line of reasoning. Indeed,

$$
\left|\gamma_{\mathrm{int}}\right|=2 \sqrt{2} \frac{\left|a_{1} a_{3}-a_{2} a_{4}\right|}{\sqrt{\alpha+\sqrt{D^{+}}}} \leq \frac{2 \sqrt{2}}{\sqrt{\alpha}}\left(\left|a_{3}-a_{4}\right| a_{1}+\left|a_{1}-a_{2}\right| a_{4}\right) \leq 4 \sqrt{2} \omega_{a}
$$

and

$$
\begin{align*}
\left|\gamma_{\mathrm{ext}}\right| & =\frac{\sqrt{2 D^{-}}}{\sqrt{\alpha+\sqrt{D^{-}}}+\sqrt{\alpha-\sqrt{D^{-}}}} \\
& \leq \sqrt{\frac{2 D^{-}}{\alpha}} \\
& =\sqrt{2} \sqrt{\left(a_{1}-a_{3}\right)^{2}+\left(a_{2}-a_{4}\right)^{2}} \sqrt{\frac{\left(a_{1}+a_{3}\right)^{2}+\left(a_{2}+a_{4}\right)^{2}}{\alpha}}  \tag{3.5}\\
& \leq 2 \sqrt{2} \max \left(\left|a_{1}-a_{3}\right|,\left|a_{2}-a_{4}\right|\right) \\
& \leq 2 \sqrt{2} \omega_{a}
\end{align*}
$$

## 4. Proof of the main result: lower bound

We suggest here a proof of the lower bound (1.7), alternative to one in [8, Theorem 3.2]. It has nothing to do with periodicity and applies to arbitrary bounded Jacobi matrices. The argument is based on some general facts from the perturbation theory (perturbation of the spectra).

We say that $\left(\lambda, \lambda^{\prime}\right)$ is a gap in the spectrum of a bounded, self-adjoint operator $T$ on a Hilbert space, if

$$
\left(\lambda, \lambda^{\prime}\right) \bigcap \sigma(T)=\emptyset, \quad \lambda, \lambda^{\prime} \in \sigma(T)
$$

The result below is likely to be well known. We provide its proof for the sake of completeness.

Lemma 4.1. Let $T_{0}$ be a bounded, self-adjoint operator on a Hilbert space with the spectrum $\sigma\left(T_{0}\right)=\left[u_{0}, v_{0}\right]$, a single interval. Let $T$ be a bounded, self-adjoint operator so that $\left\|T-T_{0}\right\| \leq \delta$. Then the length of each gap $\left(\mu^{-}, \mu^{+}\right)$in $\sigma(T)$ does not exceed $2 \delta$.

Proof. Denote by $[u, v]$ the convex hall of $\sigma(T)$, i.e., the least closed interval, which contains $\sigma(T)$. By the hypothesis on the perturbation, we have

$$
u_{0} \leq u+\delta, \quad v_{0} \geq v-\delta
$$

Assume, on the contrary, that $\mu^{+}-\mu^{-}>2 \delta$. Take the middle point of the gap $w=\left(\mu^{-}+\mu^{+}\right) / 2$, so

$$
\begin{aligned}
& w-\mu^{-}>\delta \Longrightarrow w-u>\delta, \quad w>u+\delta \geq u_{0} \\
& \mu^{+}-w>\delta \Longrightarrow v-w>\delta, \quad w<v-\delta \leq v_{0}
\end{aligned}
$$

and therefore, $w \in\left[u_{0}, v_{0}\right]=\sigma\left(T_{0}\right)$.

Pick a positive number $\tau$ so that

$$
\begin{equation*}
0<\tau<\frac{\mu^{+}-\mu^{-}-2 \delta}{2 \delta}, \quad \frac{\mu^{+}-\mu^{-}}{2}>(1+\tau) \delta \tag{4.1}
\end{equation*}
$$

We invoke the resolvent operators $R\left(z, T_{0}\right)=\left(T_{0}-z\right)^{-1}, R(z, T)=(T-z)^{-1}$, and write the equalities

$$
\begin{aligned}
R(z, T)-R\left(z, T_{0}\right) & =-R(z, T)\left(T-T_{0}\right) R\left(z, T_{0}\right) \\
R(z, T) & =\left(I-R(z, T)\left(T-T_{0}\right)\right) R\left(z, T_{0}\right)
\end{aligned}
$$

Put $z=w+i \varepsilon, \varepsilon>0$ so, by (4.1),

$$
\begin{aligned}
\|R(z, T)\|^{-1} & =\operatorname{dist}(z, \sigma(T)) \geq \sqrt{(1+\tau)^{2} \delta^{2}+\varepsilon^{2}}>(1+\tau) \delta \\
\|R(z, T)\| & <\frac{1}{(1+\tau) \delta}
\end{aligned}
$$

Then

$$
\left\|R(z, T)\left(T-T_{0}\right)\right\|<\frac{1}{1+\tau}
$$

and the operator $I-R(z, T)\left(T-T_{0}\right)$ is invertible with

$$
\left\|\left(I-R(z, T)\left(T-T_{0}\right)\right)^{-1}\right\| \leq \frac{1+\tau}{\tau}
$$

Consequently,

$$
\begin{aligned}
R\left(z, T_{0}\right) & =\left(I-R(z, T)\left(T-T_{0}\right)\right)^{-1} R(z, T) \\
\left\|R\left(z, T_{0}\right)\right\| & \leq\left\|\left(I-R(z, T)\left(T-T_{0}\right)\right)^{-1}\right\|\|R(z, T)\| \leq \frac{1}{\tau \delta}
\end{aligned}
$$

On the other hand, $w \in \sigma\left(T_{0}\right)$ implies

$$
\left\|R\left(w+i \varepsilon, T_{0}\right)\right\|=\varepsilon^{-1} \rightarrow+\infty, \quad \varepsilon \rightarrow 0+
$$

The contradiction completes the proof.

Proposition 4.2. Let $J$ be a bounded Jacobi matrix (1.1), and ( $\mu^{-}, \mu^{+}$) be a gap in its spectrum. Then

$$
\mu^{+}-\mu^{-} \leq 2\left(\omega_{b}+2 \omega_{a}\right)
$$

Proof. We apply Lemma 4.1 for the case $T=J, T_{0}=J_{\alpha, \beta}$, a constant Jacobi matrix with $\alpha$ 's along the main diagonal and $\beta$ 's along the off-diagonals, where $\alpha$ and $\beta$ are suitable constants. Then

$$
J-J_{\alpha, \beta}=\left[\begin{array}{cccccc}
\ddots & \ddots & \ddots & & & \\
& a_{-1}-\alpha & b_{0}-\beta & a_{0}-\alpha & & \\
& & a_{0}-\alpha & b_{1}-\beta & a_{1}-a & \\
& & & a_{1}-\alpha & b_{2}-\beta & a_{2}-\alpha \\
& & & \ddots & \ddots & \ddots
\end{array}\right]
$$

and, as is well known (see, e.g., [11, formula (1.3.29)]),

$$
\left\|J-J_{\alpha, \beta}\right\| \leq \sup _{j}\left|b_{j}-\beta\right|+2 \sup _{j}\left|a_{j}-\alpha\right| .
$$

Take $\inf _{j} b_{j} \leq \beta \leq \sup _{j} b_{j}, \inf _{j} a_{j} \leq \alpha \leq \sup _{j} a_{j}$, so for each $k \in \mathbb{Z}$

$$
\begin{aligned}
& \left|b_{k}-\beta\right| \leq \max \left(b_{k}-\inf _{j} b_{j}, \sup _{j} b_{j}-b_{k}\right) \leq \omega_{b} \\
& \left|a_{k}-\alpha\right| \leq \max \left(a_{k}-\inf _{j} a_{j}, \sup _{j} a_{j}-a_{k}\right) \leq \omega_{a}
\end{aligned}
$$

It follows that $\left\|J-J_{\alpha, \beta}\right\| \leq \omega_{b}+2 \omega_{a}$, and application of Lemma 4.1 completes the proof.

The lower bound (1.7) is a straightforward consequence of the latter result.
Remark 4.3. The Borg-Hochstadt theorem is known to hold for a wider class of reflectionless Jacobi matrices [10, Corollary 8.6]. One may conjecture that its quantitative form, Theorem 1.1, remains valid in this setting (at least in the finitegap case) as well.

Furthermore, the following extended version of the Borg-Hochstadt theorem is known to hold, see [4], [11, Theorem 5.13.9].

Let $J$ be a periodic Jacobi matrix (1.1) of period $p=k q$ with $k$ and $q$ integral. Suppose that all the gaps $\gamma_{j}$ are closed for $j \neq k, 2 k, \ldots,(q-1) k$. Then $J$ is periodic at period $q$.

It is quite conceivable that this version of the Borg-Hochstadt theorem is stable in the above sense. Precisely, one may conjecture that

$$
\begin{aligned}
c_{1} \max \left\{\left|\gamma_{j}\right|: j \not \equiv 0(\bmod k)\right\} & \leq \max _{|i-j|=q}\left(a_{i}-a_{j}\right)+\max _{|i-j|=q}\left(b_{i}-b_{j}\right) \\
& \leq c_{2} \max \left\{\left|\gamma_{j}\right|: j \not \equiv 0(\bmod k)\right\}
\end{aligned}
$$

holds. We observed such phenomenon for $p=4=2 \cdot 2$ in (3.2) and (3.5).

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