

# NONLINEARITY EFFECT OF HIGHLY ELASTIC DEFORMATION OF AN AMORPHOUS POLYMER

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The analysis of the kinetics of the tensile deformation of *amorphous polyimide* samples with a constant rate of  $7 \cdot 10^{-5} < \dot{\epsilon} < 6 \cdot 10^{-3} \text{ s}^{-1}$  at constant values of the temperature  $T$  in the warm state ( $T = 293 \text{ K}$ ) and under the conditions of moderate ( $T = 77 \text{ K}$ ) and deep ( $T = 4.2 \text{ K}$ ) cooling begun in [1] was continued. The three-stage deformation diagrams were registered in [1] and are shown schematically in Fig. 1, here  $\epsilon = \dot{\epsilon} t$  is the relative deformation,  $\sigma$  is the deforming stress. The initial stages of linear elasticity  $\sigma_I = M_e \epsilon$  with Young's modulus  $M_e = M_e(T)$  were recorded within  $\epsilon < 0.1$ . The next relaxation stage  $\sigma_{II}$  was interpreted in [1, 2] as a consequence of the thermomechanical activation of specific *elaston excitations* on molecular chains in an amorphous polymer, i.e., elementary acts of highly elastic deformation under the action of effective stress  $\sigma^* = \sigma - \sigma_i(\epsilon; T)$  (Fig. 2). Here,  $\sigma_i$  is the entropy component of internal stresses caused by changes in conformational entropy during the straightening of molecular chains. The stage of highly elastic deformation  $\sigma_{III}$  is also determined by the kinetics of elaston excitations, it always has an extended linear component  $\sigma_{III} = \sigma_{fe} + M_{he} \epsilon$  with the modulus of entropy elasticity  $M_{he} = M_{he}(T)$  and the conditional limit of elasticity  $\sigma_{fe}$ . But when the part of the samples were deformed in a warm state of  $T = 293 \text{ K}$ , a significant deviation from the linear law was found within the strength limit of polymer  $\epsilon < \epsilon_r$ .

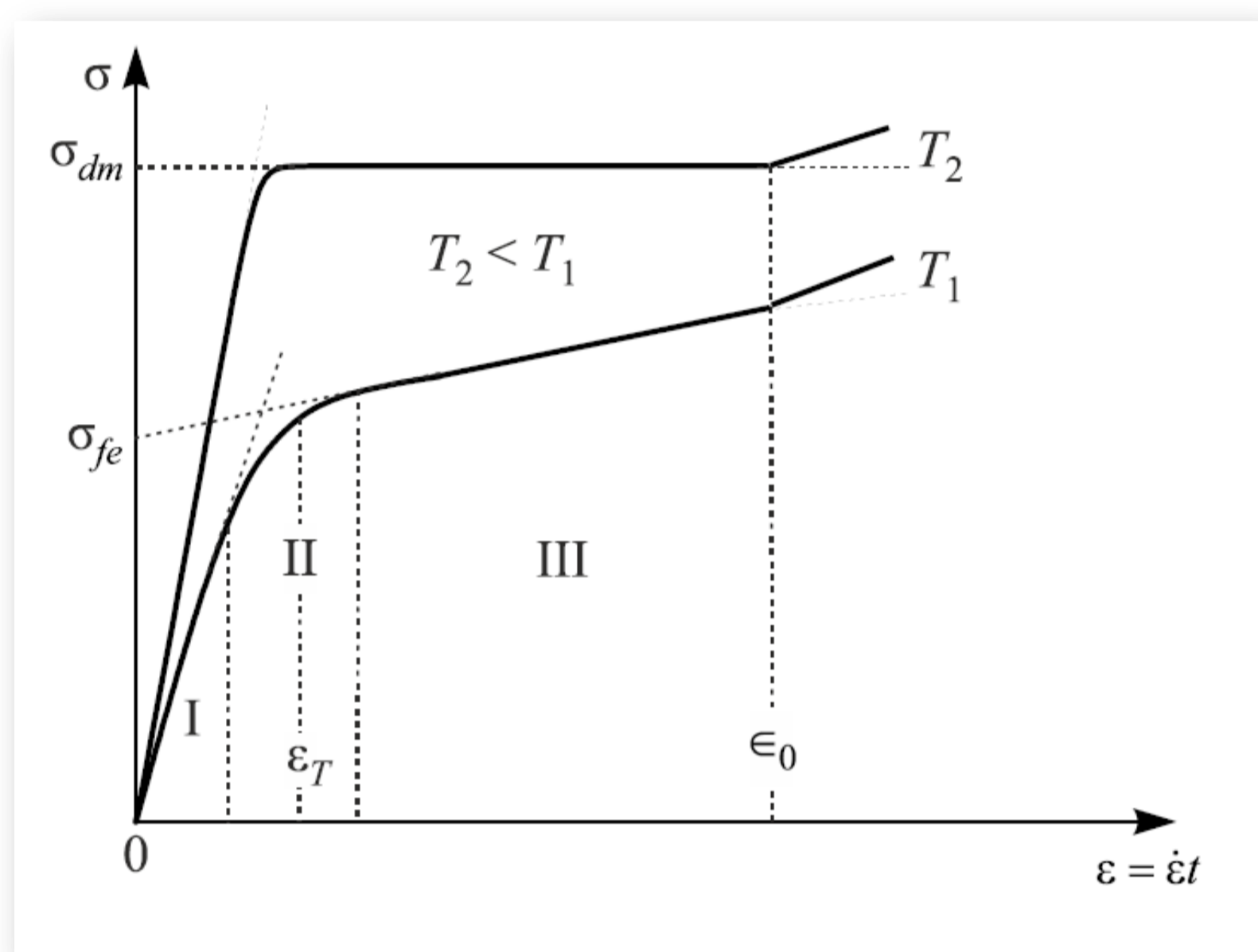


Fig. 1. General view of the deformation diagram and its transformation at polymer cooling. The lower curve  $T_1$  and  $\sigma_{fe}$  is the deformation state of a warm or moderately cooled polymer; the upper curve  $T_2$  and  $\sigma_{dm}$  is the limiting state of low-temperature deformation melting.

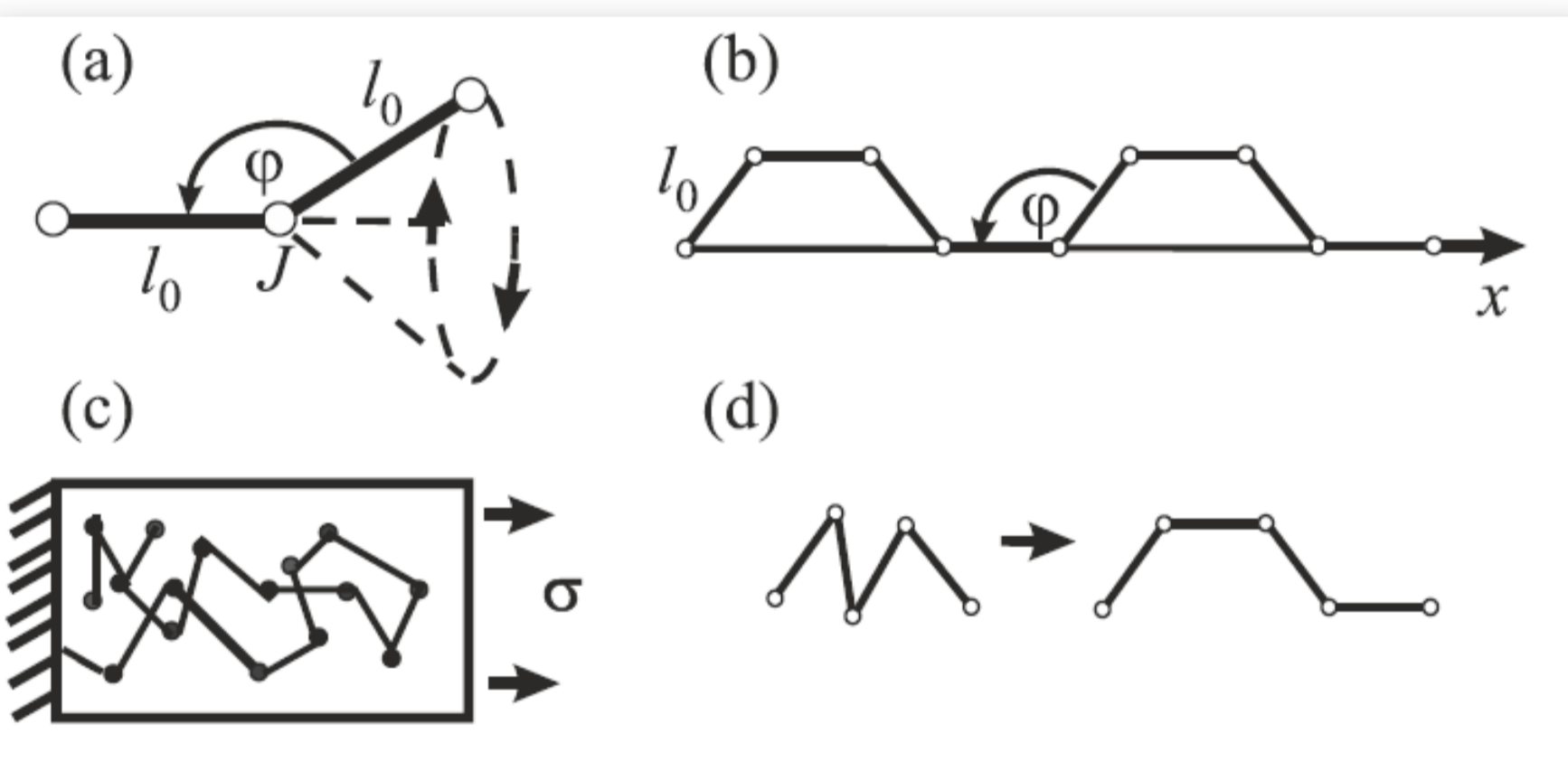


Fig. 2. Mechanical model (rods – joints) of the monomer structure of a molecular chain, a fragment of a molecular coil (sample) and an elementary act of highly elastic deformation: limited rotation of adjacent monomer units (a), molecular chain straightened along the  $x$  axis (b), conformations of molecular chains in a polymer sample (c), mechanical activation of the molecular segment (elaston excitation) (d),  $l_0$  is monomer unit (rod),  $J$  is simple covalent bond (joint),  $\phi$  is bond angle.

Table. Optimal values of rheological characteristics of polyimide samples obtained as a result of analytical approximation of their tensile test diagrams by the function (3).

$\dot{\epsilon}, \text{ s}^{-1}$	$\beta$	$\epsilon_T$	$M_e,$ MPa	$M_{he},$ MPa	$\mu_{he},$ MPa	$\sigma_{he},$ MPa	$\sigma(\epsilon_r; T, \dot{\epsilon}),$ MPa
$7 \cdot 10^{-4}$	0.6	0.05	1413	52	34	92	187
$6 \cdot 10^{-3}$	1.2	0.05	1745	65	51	100	198

The nonlinearity of the stage  $\sigma_{III}$  is interpreted as a consequence of the nonlinear dependence of  $\sigma_i$  on deformation with the second-order elastic modulus  $\mu_{he} = \mu_{he}(T)$ :

$$\sigma_i(\epsilon; T) = M_{he} \epsilon + \mu_{he} \epsilon^2. \quad (1)$$

This assumption transforms the four-parameter nonlinear deformation equation [2] into a more complex five-parameter nonlinear equation, in which the kinetics of elaston activation is complicated by the consideration of nonlinearity (1):

$$\frac{d}{d\epsilon} \sigma = M_e - (M_e - M_{he}) \beta \operatorname{sh} \left[ \frac{\sigma - M_{he} \epsilon - \mu_{he} \epsilon^2}{(M_e - M_{he}) \epsilon_T \sqrt{1 + \beta^2}} \right]. \quad (2)$$

Parameters of this equation are three isothermal coefficients of elasticity  $M_e(T)$ ,  $M_{he}(T)$ ,  $\mu_{he}(T)$  and two kinetic coefficients  $\beta$ ,  $\epsilon_T$ . Its solution with the initial condition is described by the formula:

$$\sigma(\epsilon; T, \dot{\epsilon}) = (M_e - M_{he}) \epsilon_T \sqrt{1 + \beta^2} \ln \left[ \frac{\sqrt{1 + \beta^2} + \beta + 1 + (\sqrt{1 + \beta^2} - \beta - 1) \delta(\epsilon)}{\sqrt{1 + \beta^2} + \beta - 1 + (\sqrt{1 + \beta^2} - \beta + 1) \delta(\epsilon)} \right] + 2\mu_{he} \epsilon_T^2 [1 - \delta(\epsilon)] + (M_{he} - 2\mu_{he} \epsilon_T) \epsilon + \mu_{he} \epsilon^2. \quad (3)$$

The exponential function  $\delta(\epsilon) = \exp(-\epsilon/\epsilon_T)$  determines the shape of the deformation diagram within the relaxation stage  $\sigma_{II}$ .

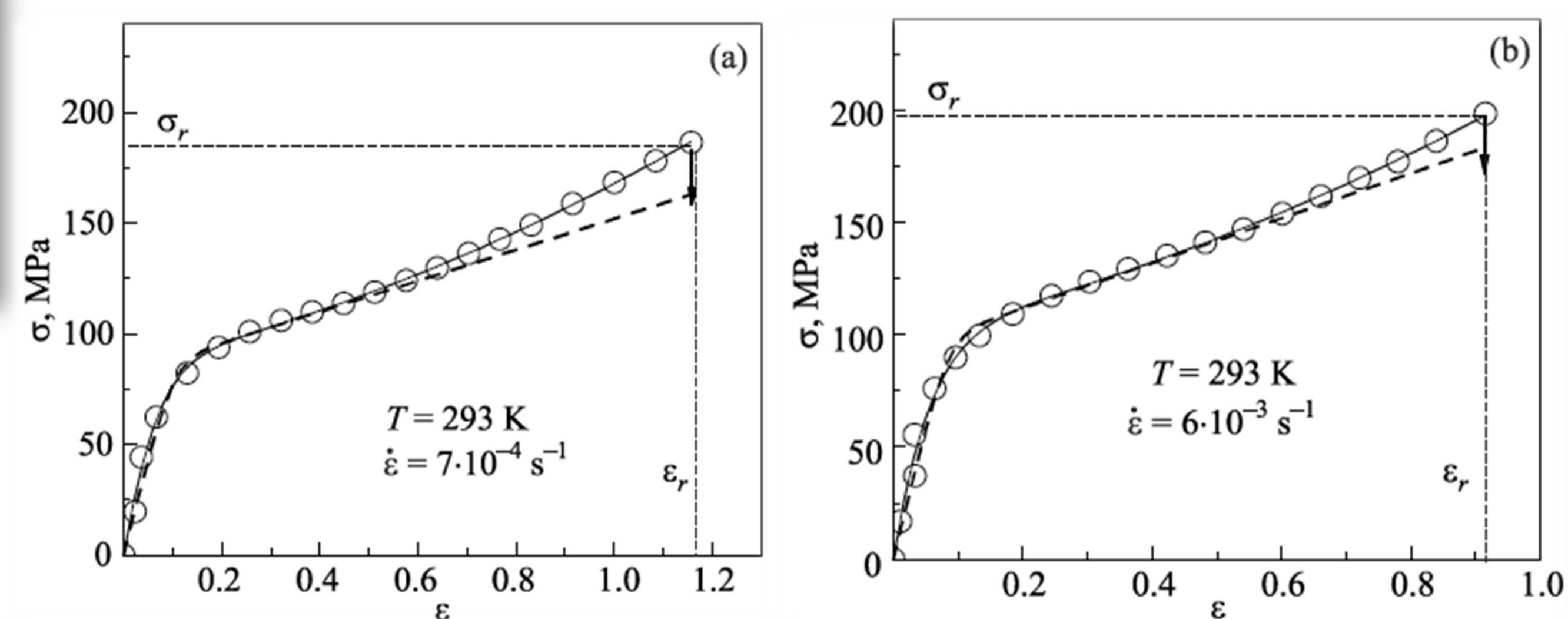


Fig. 3. The deformation diagrams of two samples of amorphous polyimide at  $T = 293 \text{ K}$  and two values of the deformation rate recorded in the experiments [1]:  $\circ$  – experimental points. Their analytical approximations are also shown: by the function  $\sigma(\epsilon)$  – formula (3), solid line; by the function  $\sigma(\epsilon)$  without taking into account nonlinearity, dash line.

- [1] V. D. Natsik, H. V. Rusakova, S. V. Lubenets, V. A. Lototskaya, and L. F. Yakovenko, *Fiz. Nizk. Temp.* 49 569 (2023) [*Low Temp. Phys.* 49, 521 (2023)], <https://doi.org/10.1063/10.0017812>.  
[2] V. D. Natsik, H. V. Rusakova, *Low Temp. Phys.* 49, 228 (2023). <http://dx.doi.org/10.1063/10.0016875>.